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<u>Example: Solutions</u> <u>Involving Inhomogeneous</u> <u>Dielectrics</u>

We again consider the parallel plate problem, only this time the plates are filled with two **different** layers of dielectric !

 $\begin{array}{c} z = -(d_1 + d_2) \\ z = -d_2 \\ \hline \\ z = 0 \\ \hline \\ z \end{array}$

The top dielectric layer has thickness d_1 and permittivity ε_1 , while the bottom layer has thickness d_2 and permittivity ε_2 .

Q: How do we determine the fields within these dielectric layers?

A: Begin by defining the electric potential function in **each** dielectric:

$$V_1(\overline{r}) \doteq$$
 electric potential function in material ε_1

$$V_2(\bar{r}) \doteq$$
 electric potential function in material ε_2

Note as before, these functions will be **independent** of coordinates x and y. Therefore:

$$V_1(\overline{r}) = V_1(z)$$
 and $V_2(\overline{r}) = V_2(z)$

Each of these functions must satisfy **Laplace's equation**, therefore:

$$\nabla^2 V_1(\overline{r}) = \frac{\partial^2 V_1(z)}{\partial z^2} = 0$$

$$\nabla^{2}V_{2}(\overline{r}) = \frac{\partial^{2}V_{2}(z)}{\partial z^{2}} = 0$$

We **know** from an earlier handout that these solutions will have the form:

$$V_1(z) = \mathcal{C}_{1a} z + \mathcal{C}_{1b}$$

$$V_2(z) = \mathcal{C}_{2a} z + \mathcal{C}_{2b}$$

YIKES! We now have FOUR unknowns $(C_{1a}, C_{1b}, C_{2a}, C_{2b})!$ Therefore, we need **four boundary conditions** to determine these constants.

We know for starters, that :

$$V_1(z = -(d_1' + d_2')) = V_0 = -C_{1a}(d_1' + d_2') + C_{1b}$$

$$V_2(z=0)=0=C_{2a}(0)+C_{2b}$$

Therefore, it is evident that $C_{2b} = 0$. But, what about the other **three** constants?

We need two more boundary conditions!

Another boundary condition can be found at the interface between the two dielectrics (i.e., at $z = -d_2$). The electric pontential at this boundary must be **the same** for **both** expressions:

$$V(z = -d_2) = V_1(z = -d_2) = V_2(z = -d_2)$$

As a result, we get the boundary equation:

$$V_{1}(z = -d_{2}) = V_{2}(z = -d_{2})$$
$$C_{1a}(-d_{2}) + C_{1b} = C_{2a}(-d_{2})$$

where we have used the fact that $C_{2b} = 0$.

But, we still need one more boundary condition.

Recall a boundary condition for dielectric interfaces is:

$$\mathbf{D}_{1n}\left(\overline{\mathbf{r}}\right) = \mathbf{D}_{2n}\left(\overline{\mathbf{r}}\right)$$

We can apply this boundary condition as well! Recall:

$$\mathsf{D}(\overline{\mathsf{r}}) = \varepsilon \, \mathsf{E}(\overline{\mathsf{r}}) = -\varepsilon \nabla \, \mathcal{V}(\overline{\mathsf{r}})$$

where for this case:

$$\nabla V_1(\overline{\mathbf{r}}) = \frac{\partial V_1(z)}{\partial z} \hat{a}_z = \mathcal{C}_{1a} \hat{a}_z$$

and

$$\nabla V_{2}(\overline{\mathbf{r}}) = \frac{\partial V_{2}(z)}{\partial z} \hat{a}_{z} = C_{2a} \hat{a}_{z}$$

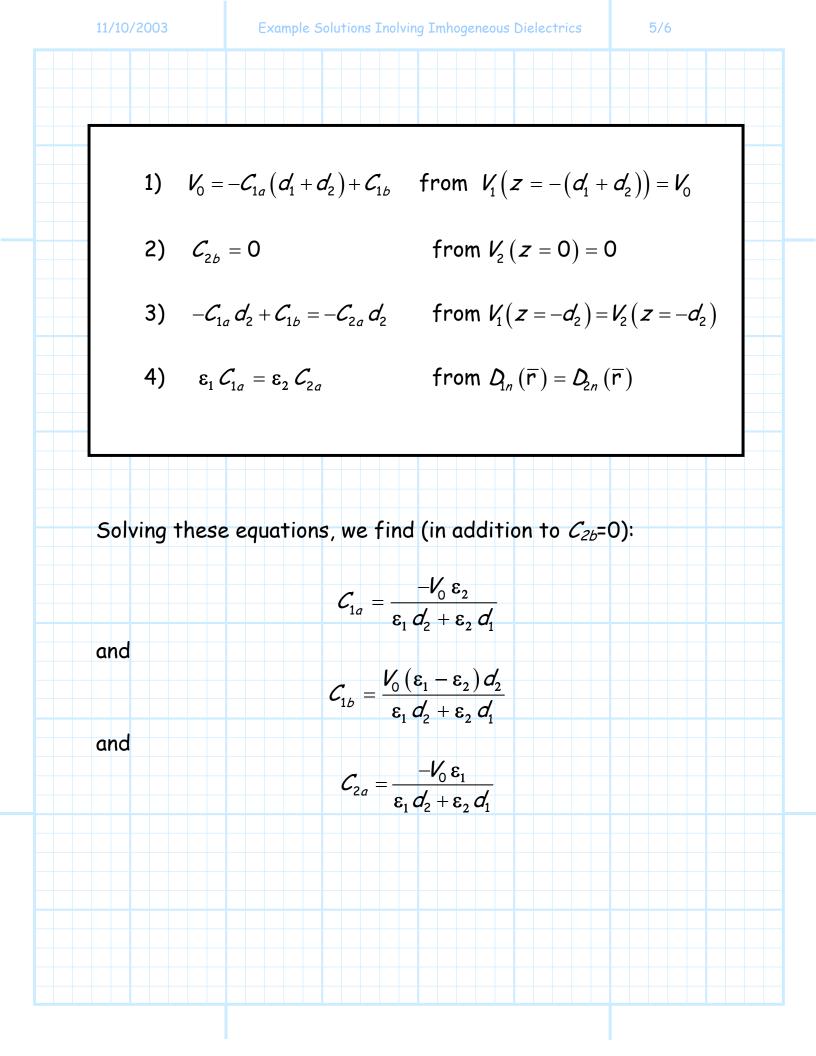
Therefore:

$$\mathbf{D}_{1}(\mathbf{\overline{r}}) = -\varepsilon_{1} \mathcal{C}_{1a} \hat{a}_{z}$$
$$\mathbf{D}_{2}(\mathbf{\overline{r}}) = -\varepsilon_{2} \mathcal{C}_{2a} \hat{a}_{z}$$

Since $\hat{a}_n = \hat{a}_z$, we find from this boundary condition that:

$$\mathcal{D}_{1n}\left(\overline{\mathbf{r}}\right) = \mathcal{D}_{2n}\left(\overline{\mathbf{r}}\right)$$
$$-\varepsilon_{1} \mathcal{C}_{1a} = -\varepsilon_{2} \mathcal{C}_{2a}$$

Now we can **solve** for our constants ! Recall our **four** equations are:



Inserting these results into the equations:

$$V_1(z) = C_{1a} z + C_{1b}$$

$$V_2(z) = \mathcal{C}_{2a} z + \mathcal{C}_{2b}$$

provides the electric potential field in each region within the parallel plates. From this result, we can determine the **electric field** and the **electric flux density** within each region, as well as the charge density on each plate.